Soft Contribution to the Hard Ridge

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CATHIE/TECHQM, December 14-18, 2009

arXiv:0910.3590 [nucl-th]

GM, Sean Gavin

Phys.Rev.C79,051902, arXiv:0806.4718 [nucl-th]

Sean Gavin, Larry McLerran, G. M.

The Ridge

- Hard Ridge: jet trigger
- Soft Ridge: no trigger
- Flow and jets

Long Range Correlations

- PHOBOS Data
- Flux Tubes, Glasma, and Correlations

Comparison to Experiment

- Blast Wave Flow + Glasma
- p_t Dependence
- Soft and Hard Ridge from STAR

Hard Ridge: Jet + Associated Particles

Measure

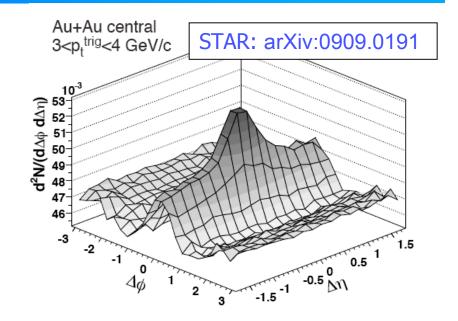
- High pt trigger
- Yield of associated particles per trigger

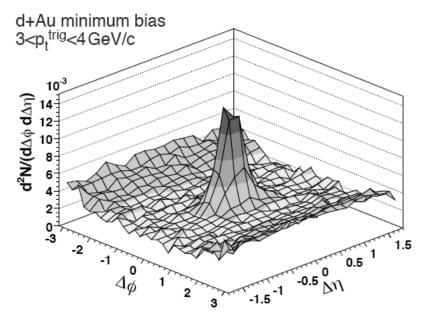
$$\frac{1}{N_{trig}} \frac{d^2 N}{d\Delta \phi \ d\Delta \eta}$$

Hard Ridge: Near Side Peak

- Peaked near $\Delta \phi = 0$
- Broad in Δη

How does the formation of the ridge at large $\Delta \eta$ depend on jets?





Soft Ridge: Untriggered Correlations

two particle correlations with no jet tag

STAR: arXiv:0806.2121

Measure

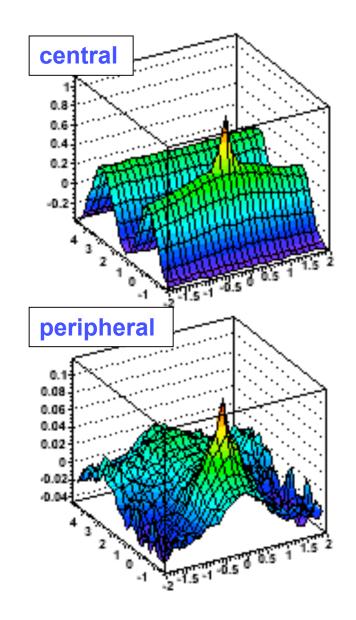
$$\frac{\Delta \rho(\eta, \phi)}{\sqrt{\rho_{ref}}} = \frac{\text{pairs} - (\text{singles})^2}{\text{singles}}$$

Soft and Hard Ridges Similar

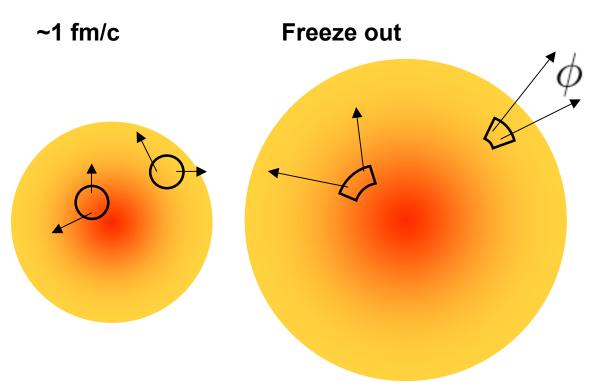
- Peaked near $\Delta \phi = 0$
- Wider in Δφ than hard ridge
- Broad in Δη
- Jet peak?

Common Features

- Δη width increases with centrality
- peripheral ~ proton+proton



Near Side ϕ Peak: Flow



Azimuthal correlations come from flow.

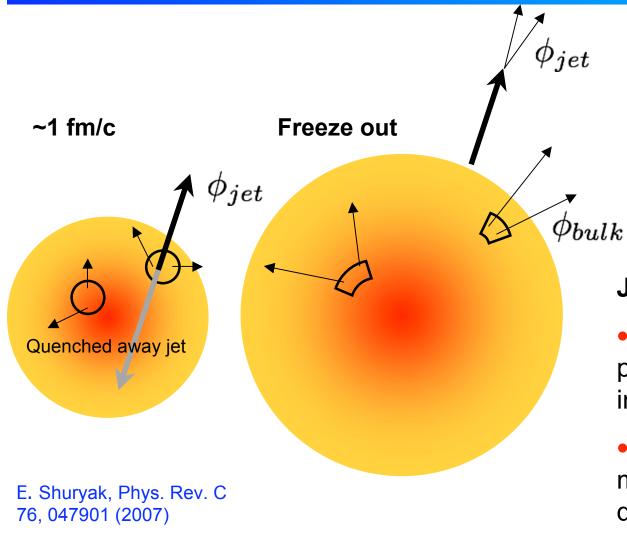
- Particles are pushed to higher p_t and and focused to a smaller azimuthal angle depending on the push.
- The ridge should narrow in
 φ with increasing p_t cuts.

- Mean flow depends on position
- Opening angle for each fluid element depends on radial position

$$\vec{v}_t \sim \lambda \vec{r}_t$$

$$\phi \sim v_{th}/v_t \propto (\lambda r_t)^{-1}$$

Near Side φ Peak: Jets + Flow



Claim:

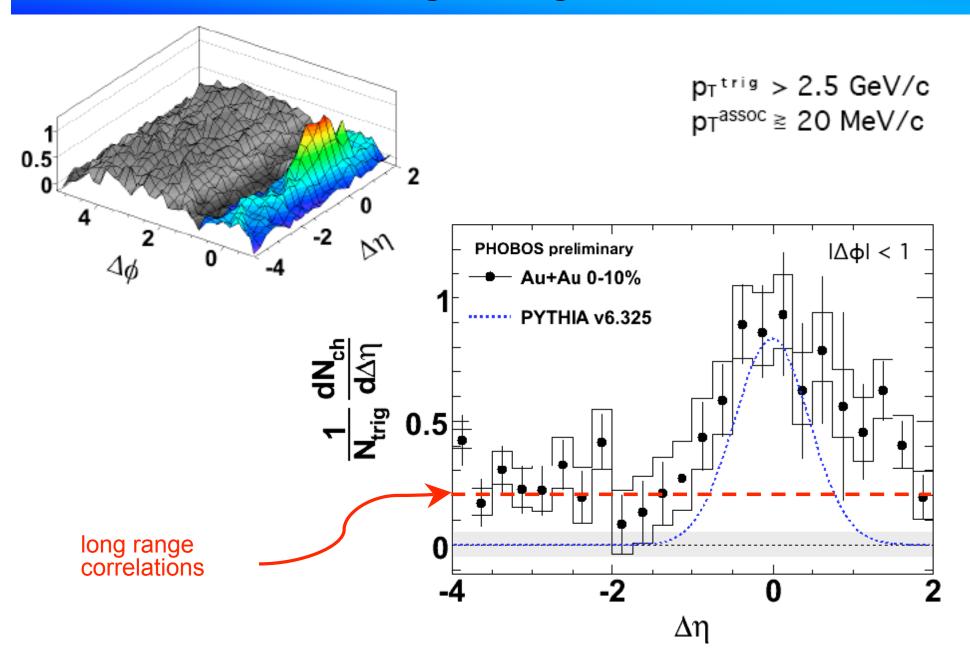
Soft ridge explained by bulk flow

Hard ridge: additional jet-bulk contribution

Jet Correlations With Bulk

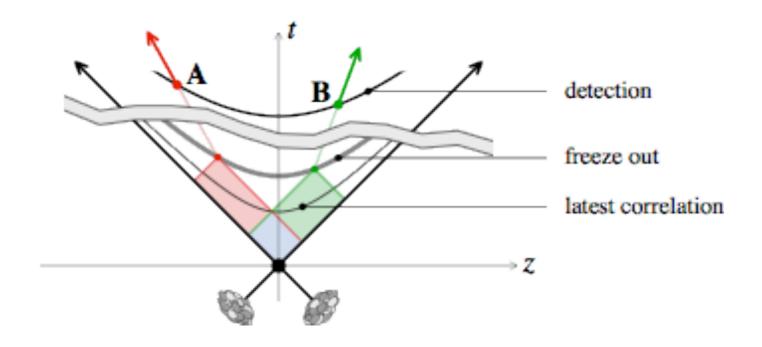
- Correlation of flow and jet particles if produced nearby in transverse plane
- Surviving jets tend to be more radial, due to quenching.
- Bulk particles are pushed into the radial direction by flow

PHOBOS: Long Range Correlations



Why Long Range Correlations?

Dumitru, Gelis, McLerran, Venugopalan, arXiv:0804.3858

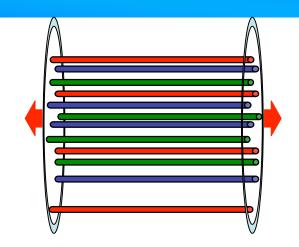


- must originate at the earliest stages of the collision
- like super-horizon fluctuations in the Universe
- information on particle production mechanism

Flux Tubes and Glasma

Flux Tubes: longitudinal fields early on





- ullet Number of flux tubes $\dfrac{Area}{area/tube} \propto Q_s^2 R_A^2$
- Tubes—quarks+gluons Single flux tube phase space density of gluons $\sim \alpha_s^{-1}(Q_s)$
- Gluon rapidity density
 Kharzeev & Nardi

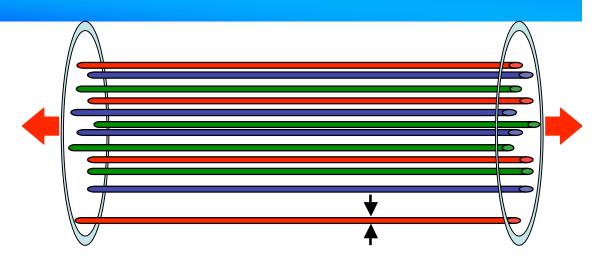
$$\frac{dN}{dy} \propto \alpha_s^{-1}(Q_s)Q_s^2 R_A^2$$

Flux Tubes and Correlations

Correlation function

- Partons from the same tube are correlated
- Rapidity reach: Causally disconnected

See: Dusling et. al. arXiv:0911.2720



flux tube transverse

$$c(\mathbf{x}_1, \mathbf{x}_2) \sim \mathcal{R} \, \delta(\vec{r}_{t,1} - \vec{r}_{t,2})$$

- size $\sim Q_s^{-1} << R_A$
- Correlation Strength $\mathcal{R} \propto \langle \#tubes \rangle^{-1} = (Q_s R_A)^{-2}$
- Long range Glasma fluctuations scale the phase space density

$$\mathcal{R}\frac{dN}{dy} \propto \alpha_s^{-1}(Q_s^2)$$

Dumitru, Gelis, McLerran & Venugopalan; Gavin, McLerran & GM

Energy and centrality dependence

$$Q_s^2 \propto (\sqrt{s})^{1/3} (N_{part})^{1/3}$$

Energy and System Dependence

$$\Delta \rho \equiv pairs - (singles)^2$$

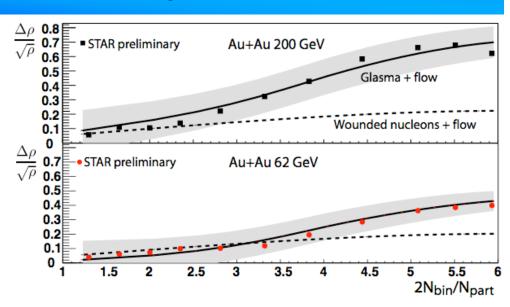
$$\propto \iint c(\vec{x}_1, \vec{x}_2) f(\vec{x}_1, \vec{p}_1) f(\vec{x}_2, \vec{p}_2)$$

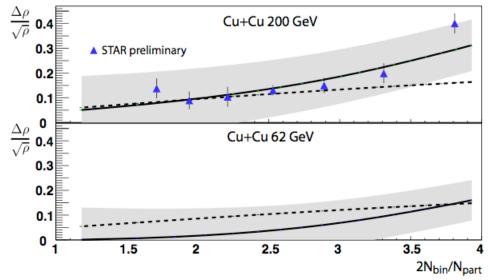
Blast Wave

- Boltzmann Dist. $\rightarrow f(p,x)$
- Scale factor to fit 200 GeV only
- Centrality dependence on blast wave parameters (v and T) → 10% uncertainty
- Blast wave only (dashed) fails

Glasma Dependence

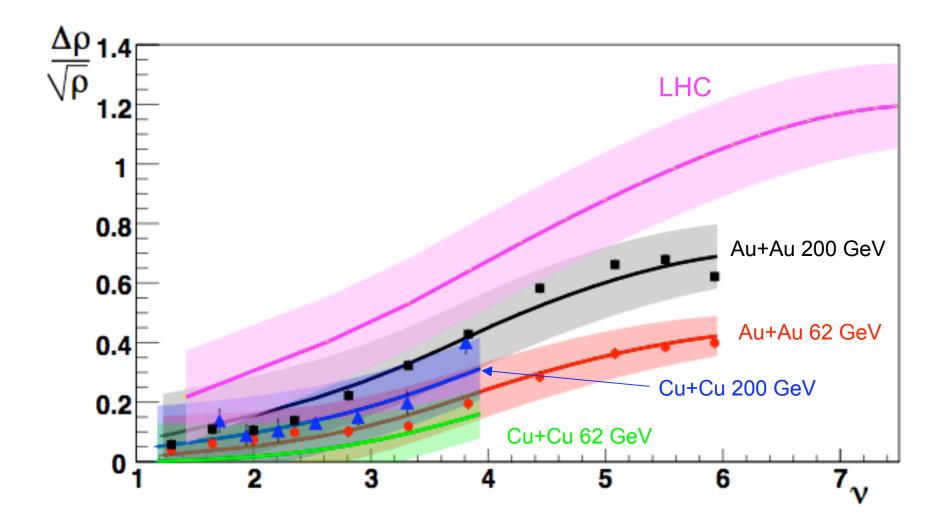
$$\frac{\Delta \rho}{\sqrt{\rho}} = \mathcal{R}\frac{dN}{dy} \times (blast\ wave)$$





• Q_s dependence: 200 GeV Au+Au \Rightarrow 62 GeV, Cu+Cu

Comparison and LHC

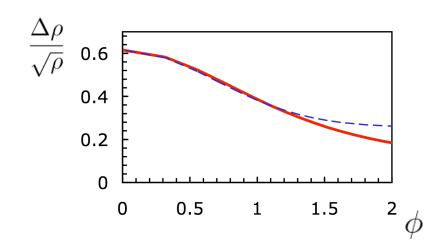


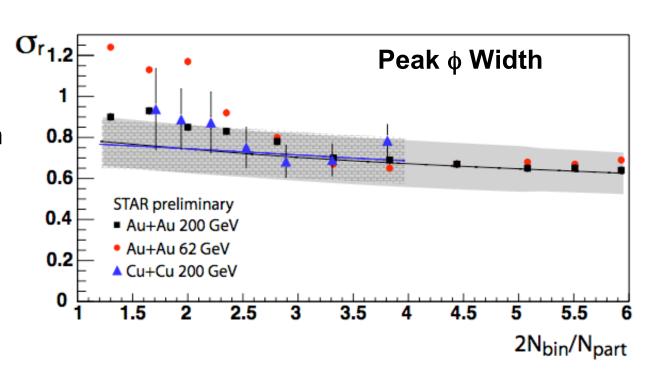
Caution: Blast Wave parameters for LHC taken from Au+Au 200 GeV

Angular Correlations

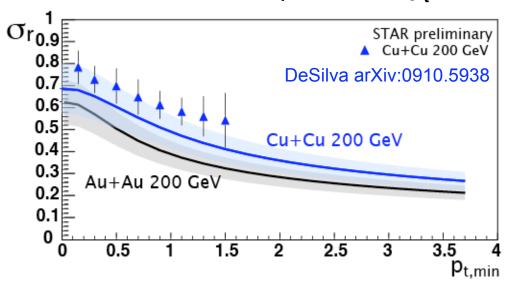
Fit using Gaussian + offset

- Range: $-\pi/2 < \phi < \pi/2$
- Error band: 20% shift in fit range
- Uncertainty due to experimental definition of peak
- Computed angular width is approximately independent of energy
- The width should decrease with increasing p_t range

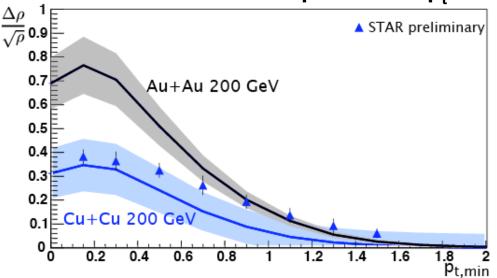




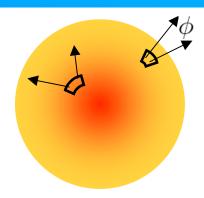
Soft Ridge vs. pt



Most Central Amplitude vs. pt



Examine bulk correlations in different p_t ranges



- The amplitude drops and the azimuthal width narrows with increasing p_{t.min}
- Bulk correlations alone might not explain the data at higher p_t
- Jet-Bulk and Jet-Jet correlations should have an increasing effect with p_t
- Jet contributions should force the correlation width to approach the jet correlation width

Jets + Glasma

Jet-Bulk correlation function

$$c(\mathbf{x}_1, \mathbf{x}_2)$$

 $\sim \mathcal{R}_{JB} \delta(\vec{r}_{t,jet} - \vec{r}_{t,tube})$

Correlation strength

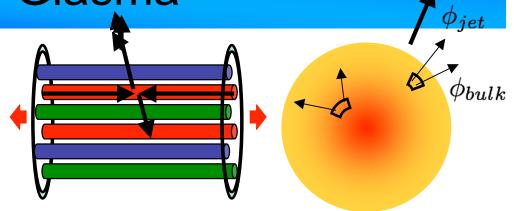
$$\mathcal{R}_{_{JB}}=\mathcal{R}$$

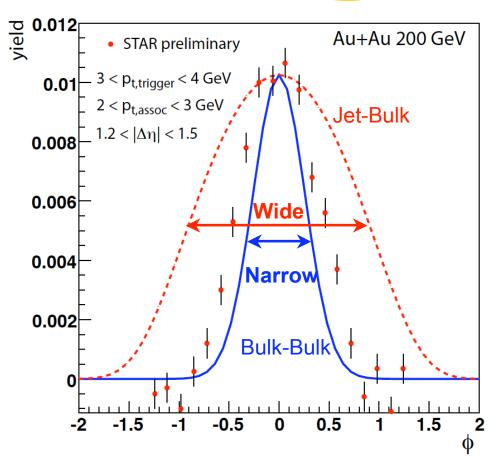
 Yield of associated particles per jet trigger; different p_t ranges

$$\frac{1}{N_{trig}}\frac{d^2N}{d\Delta\phi~d\Delta\eta}$$

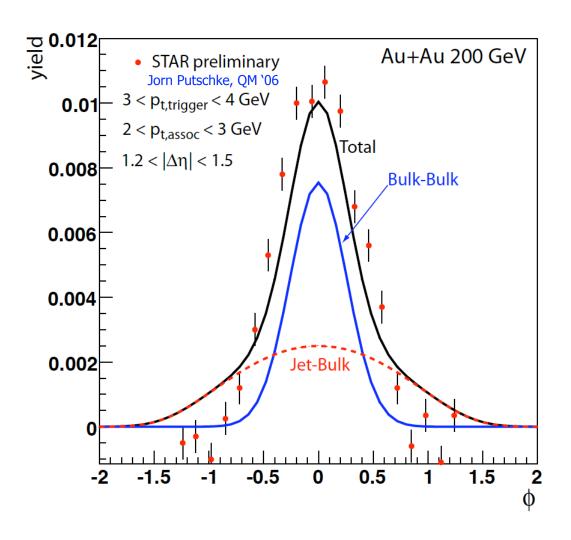
• $f(x_1, p_1) \rightarrow \text{jet p}_t \text{ range}$ $f(x_2, p_2) \rightarrow \text{bulk associated p}_t \text{ range}$

Jet-Bulk width similar to E. Shuryak, Phys. Rev. C 76, 047901 (2007)





Hard Ridge



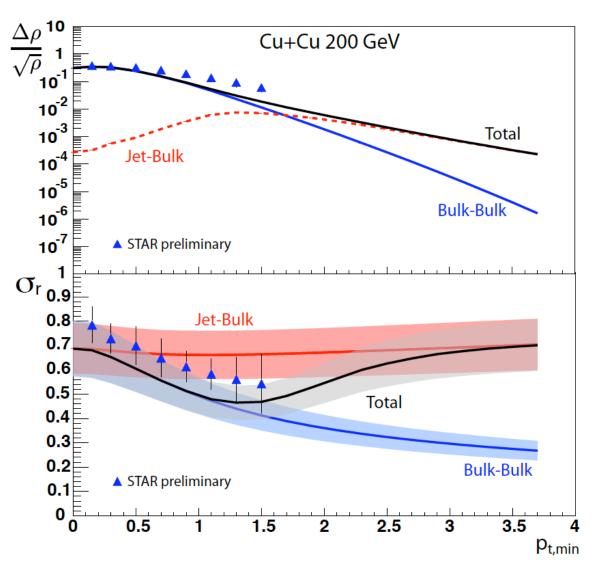
dN/dp_t constrains jet fraction

- Bulk particles: Blast Wave
- Jet particles: Total BW

Jets + Flow Fit the Hard Ridge

- Bulk-Bulk correlations ~70%.
- Bulk-Bulk + Jet-Bulk better azimuthal agreement

The Ridge: From Soft to Hard



Bulk Correlations

- •Amplidude decreases with p_{t,min}
- Narrow width from flow alone

Jet+Bulk Correlations

- Jet contribution dominates with increasing p_{t,min}
- σ_r widening at large $p_{t,min}$ would indicate significant contribution from jet correlations out in the ridge

Summary

Ridge Azimuthal Width

- Flow induces angular correlations
- Azimuthal width vs. p_t can distinguish flow from jets

Long Range Correlations

- PHOBOS measurement
- Implications on particle production mechanism

Glasma + Blast Wave

- Blast Wave fixed by single particle spectra
- Glasma fixed by dN/dy and 200 GeV Au+Au
- Predicts the height and azimuthal width of the Soft and Hard Ridge
- Predict energy, centrality, system, and p₁ dependence

Bulk Correlations Dominate the Hard Ridge

Backup Slides

Hard vs. Soft Ridge

hard ridge explanations -- jet interactions with matter

- N. Armesto, C.A. Salgado, U.A. Wiedemann, Phys. Rev. Lett. 93, 242301 (2004)
- P. Romatschke, Phys. Rev. C 75, 014901 (2007)
- A. Majumder, B. Muller, S. A. Bass, Phys. Rev. Lett. 99, 042301 (2007)
- C. B. Chiu, R. C. Hwa, Phys. Rev. C 72, 034903 (2005)
- C. Y. Wong, arXiv:0712.3282 [hep-ph]
- R. C. Hwa, C. B. Yang, arXiv:0801.2183 [nucl-th]
- T. A. Trainor, arXiv:0708.0792 [hep-ph]
- A. Dumitru, Y. Nara, B. Schenke, M. Strickland, arXiv:0710.1223 [hep-ph]
- E. V. Shuryak, Phys. Rev. C 76, 047901 (2007)
- C. Pruneau, S. Gavin, S. Voloshin, Nucl.Phys.A802:107-121,2008

soft ridge -- similar but no jet -- collective behavior

- S. Gavin and M. Abdel-Aziz, Phys. Rev. Lett. 97, 162302 (2006)
- S. A. Voloshin, Phys. Lett. B 632, 490 (2006)
- S. Gavin and G. Moschelli, arXiv:0806.4366 [nucl-th]
- A. Dumitru, F. Gelis, L. McLerran and R. Venugopalan, arXiv:0804.3858 [hep-ph]
- S. Gavin, L. McLerran, G. Moschelli, arXiv:0806.4718; arXiv:0910.3590 [nucl-th]
- F. Gelis, T. Lappi, R. Venugopalan, arXiv:0807.1306 [hep-ph]
- J. Takahashi et. al. arXiv:0902.4870 [nucl-th]

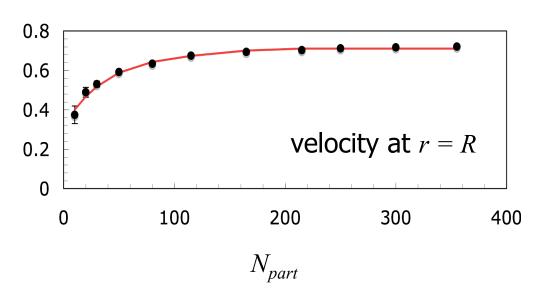
Blast Wave Single Particle Fits

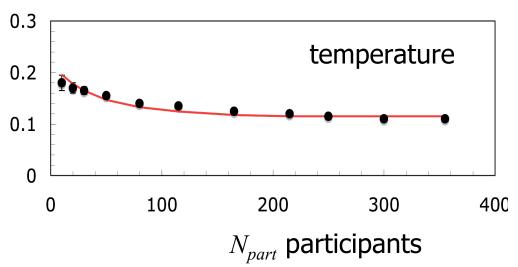
fit momentum spectra in 200 GeV Au+Au

10% systematic uncertainty in scale of v and T

62 GeV Au+Au: 5% smaller v, 10% smaller T

Akio Kiyomichi, PHENIX





Blast Wave and the Correlation Function

Schnedermann, Sollfrank & Heinz

Single Particle Spectrum

Correlation Function

$$\rho_{1}(\vec{p}) \equiv \frac{dN}{dyd^{2}p_{t}} = \int_{freezout} f(\vec{x}, \vec{p})$$

$$\gamma_{t}\vec{v}_{t} = \lambda \vec{r}$$

 $\gamma_t \dot{v_t} = \lambda r$ A Hubble like

expansion in used in a

Boltzmann Distribution

$$\Delta \rho(\vec{p}_1, \vec{p}_2) \equiv pairs - (singles)^2$$

$$\Delta \rho(\vec{p}_1, \vec{p}_2) = \iint_{freezout} c(\vec{x}_1, \vec{x}_2) f(\vec{x}_1, \vec{p}_1) f(\vec{x}_2, \vec{p}_2)$$

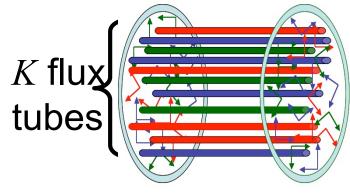
$$\Delta \rho(\eta,\phi) = \iint\limits_{momenta} \Delta \rho(\vec{p}_{\scriptscriptstyle 1},\vec{p}_{\scriptscriptstyle 2})$$

$$\rho_{ref} = \iint_{momenta} \rho_1(\vec{p}_1)\rho_1(\vec{p}_2)$$

$$\iint_{momenta} \Delta \rho = \iint_{positions} c = \langle N \rangle^2 \mathcal{R}$$

Correlation Strength

strength R



$$\langle N \rangle^{2} \mathbf{R} = \iint_{volume} c(x_{1}, x_{2}) =$$

$$= \iint_{volume} \left\{ n_{2}(x_{1}, x_{2}) - n_{1}(x_{1}) n_{1}(x_{2}) \right\} = \langle N(N-1) \rangle - \langle N \rangle^{2}$$

K flux tubes, assume

 $\langle N \rangle_K = \mu K, \qquad \langle N^2 \rangle_K - \langle N \rangle_K^2 = \sigma^2 K$

K variesevent-by-event

 $\langle N \rangle = \mu \langle K \rangle, \qquad \langle N^2 \rangle - \langle N \rangle^2 = \sigma^2 \langle K \rangle + \mu^2 (\langle K^2 \rangle - \langle K \rangle^2)$

$$\mathbf{R} = \frac{\sigma^2 - \mu}{\mu^2} \frac{1}{\langle K \rangle} + \frac{\langle K^2 \rangle - \langle K \rangle^2}{\langle K \rangle^2}$$

fluctuations per tube

number of tubes

Jet Correlation Strength

$$\mathcal{R} = \frac{\langle N(N-1)\rangle - \langle N\rangle^2}{\langle N\rangle^2}$$

$$\mathcal{R}_{\scriptscriptstyle JB} = \frac{\langle N_{\scriptscriptstyle J} N_{\scriptscriptstyle B} \rangle - \langle N_{\scriptscriptstyle J} \rangle \langle N_{\scriptscriptstyle B} \rangle}{\langle N_{\scriptscriptstyle J} \rangle \langle N_{\scriptscriptstyle B} \rangle}$$

$$=\frac{\alpha\beta\langle N(N-1)\rangle - \alpha\beta\langle N\rangle\langle N\rangle}{\alpha\beta\langle N\rangle\langle N\rangle}$$

Pruneau, Gavin, Voloshin Phys.Rev. C66 (2002) 044904

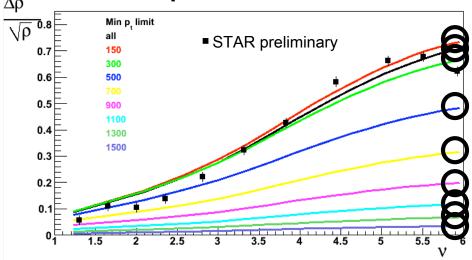
$$\begin{split} \langle N_{\scriptscriptstyle B} \rangle &= \beta \langle N \rangle \\ \langle N_{\scriptscriptstyle J} \rangle &= \alpha \langle N \rangle \\ \langle N_{\scriptscriptstyle J} N_{\scriptscriptstyle B} \rangle &= \alpha \beta \langle N(N-1) \rangle \end{split}$$

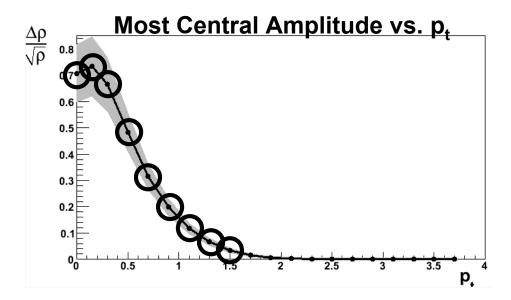


$$\mathcal{R}_{_{JB}}=\mathcal{R}$$

Soft Ridge vs. pt







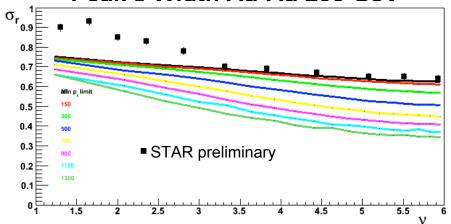
$$\left(\frac{\Delta\rho}{\sqrt{\rho_{ref}}}\right)_{p_t>p_{t,min}}$$

$$= \frac{\iint\limits_{p_{t,min}} \Delta \rho(\vec{p}_{t1}, \vec{p}_{t2})}{\{\iint\limits_{p_{t,min}} \rho_{1(\vec{p}_{t1})} \rho_{1(\vec{p}_{t2})}\}^{1/2}}$$

- Increase the lower p_t limit of the soft ridge calculation toward the hard ridge range.
- As the lower p_t limit is increased less particles are available for correlations.
- Correlation amplitude for the most central collision plotted vs. the lower p_t limit.

Soft Ridge vs. p_t

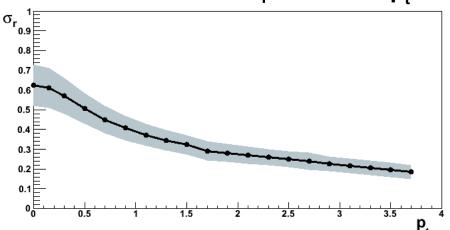
Peak & Width Au-Au 200 GeV



Angular width from

$$\left(\frac{\Delta\rho}{\sqrt{\rho_{ref}}}\right)_{p_t>p_{t,min}}$$

Most Central φ Width vs. p_t



 Higher p_t particles received a larger radial push ⇒ narrower relative angle.

Quenching + Flow

- Surviving jets tend to be more radial, due to quenching.
- Jet path

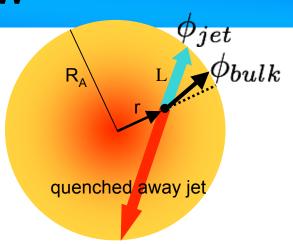
$$L(r,\phi_{jet}) = \sqrt{R_A^2 - r^2 \sin^2(\phi_{jet})} - r \cos(\phi_{jet})$$

 $S(\vec{x_1}, \vec{x_2}) = e^{\frac{-L(r, \phi_{jet})}{\rho\sigma}}$

Survival probability

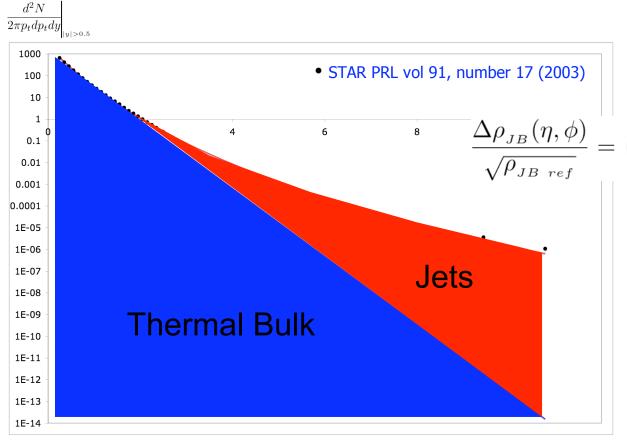
 $P_{prod}(r) \propto \left(1 - \frac{r^2}{R^2}\right)$ Production probability

 $f(\vec{x}, \vec{p}) = \frac{A}{p^n} P_{prod}(r) S(r, \phi_{jet})$ Jet Distribution



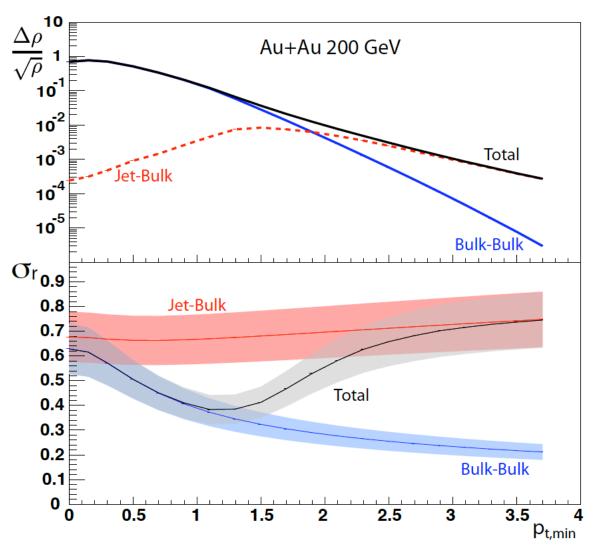
E. Shuryak, Phys. Rev. C 76, 047901 (2007)

Two Contributions



 $\frac{\Delta \rho_{JB}(\eta,\phi)}{\sqrt{\rho_{JB}\ ref}} = \begin{pmatrix} \int \Delta \rho_{JB}(\vec{p}_{t1},\vec{p}_{t2}) \\ \int Both the Bulk-Bulk and Jet-Bulk-Bulk and Jet-Bulk-Bulk and Jet-Bulk contributions are used by the total.$

The Ridge: From Soft to Hard



Bulk Correlations

- •Amplidude decreases with p_{t,min}
- Narrow width from flow alone

Jet+Bulk Correlations

- Jet contribution dominates with increasing p_{t.min}
- σ_r widening at large $p_{t,min}$ would indicate significant contribution from jet correlations out in the ridge